

哈尔滨工业大学航天学院力学系

高等弹性动力学第5讲

Biot孔隙介质的几个弹性系数

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石油勘探经常需要孔隙介质岩石模型

石油天然气存于孔隙岩石。孔、渗、饱是描述储层的基本参数。

孔隙介质中弹性波的一种传播模型——**Biot** 波动理论



著名力学家，铁木辛柯奖得主
Maurice A. Biot 1905-1985

建立了含液体孔隙岩石的静力学模型（1941）

建立了井孔声波的力学模型，最早计算了井孔声场的伪瑞利波、斯通利波的频散曲线（1952）

Biot's equations for waves in porous media (1956-1962)
建立了描述孔隙介质流体-固体耦合运动的波动力学，预测了慢纵波的存在。

Biot Conferences on Poromechanics

——致力于研究孔隙介质力学与声学的国际会议

5th Biot Conference, Vienna, 2013，胡恒山组3篇论文

4th 2009 Columbia U., USA

3rd 2005 Oklahoma, USA，胡恒山宣读论文

$$\varepsilon = a_{11}\sigma + a_{12}p$$

$$\zeta = a_{21}\sigma + a_{22}p$$



$$\varepsilon = \frac{1}{K}\sigma + a_{12}p$$

$$\zeta = a_{21}\sigma + a_{22}p$$

$$\delta\sigma \neq 0$$

$$a_{11} = \frac{1}{K} = \left. \frac{\delta\varepsilon}{\delta\sigma} \right|_p$$

$$p = 0$$

$$\frac{1}{K}$$

*Drained
Compression
Coefficient*

$$\frac{1}{K_d}$$

$$(-\sigma) \rightarrow (-\varepsilon) \uparrow$$

排水体积压缩系数

围压增大时，样品体积减小

$$\varepsilon = a_{11}\sigma + a_{12}p$$

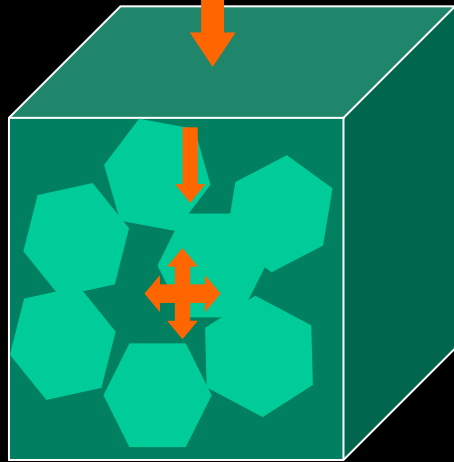
$$\zeta = a_{21}\sigma + a_{22}p$$



$$\varepsilon = \frac{1}{K}\sigma + \frac{1}{H^w}p$$

$$\zeta = a_{21}\sigma + a_{22}p$$

$$\zeta \neq 0$$



$$a_{12} = \frac{1}{H^w} = \left. \frac{\delta\varepsilon}{\delta p} \right|_{\sigma}$$

$$\sigma = 0$$

$$p \uparrow \rightarrow \varepsilon \uparrow$$

$$\frac{1}{H^w}$$

*poroelastic
expansion
coefficient*

like Heat expansion

围压不变，
增加孔压 (注入流体)，则样品膨胀
孔隙弹性膨胀系数

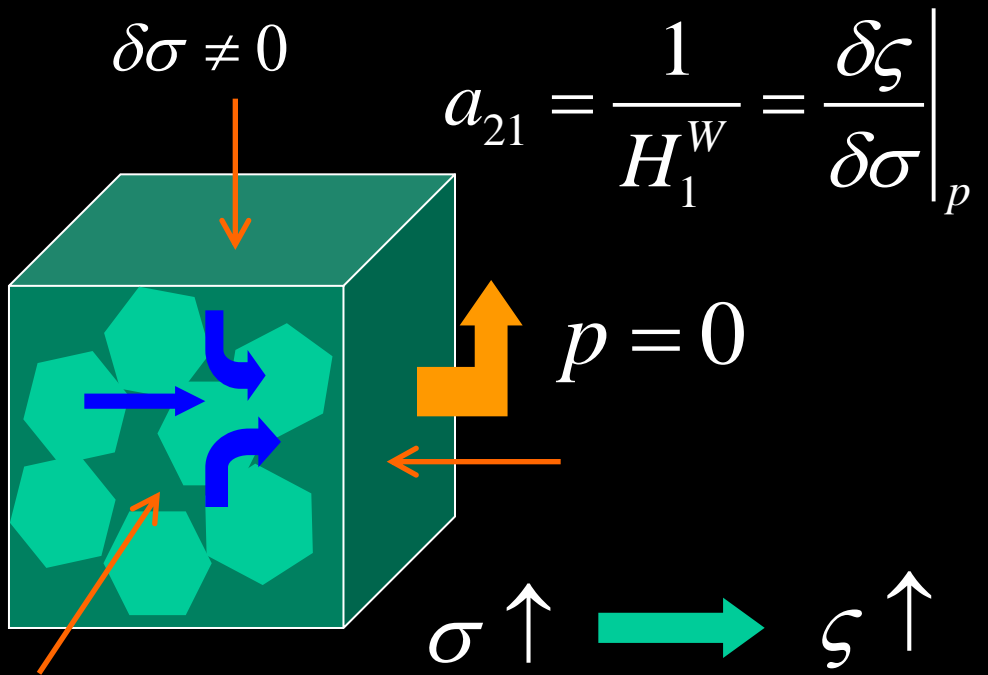
$$\varepsilon = a_{11}\sigma + a_{12}p$$

$$\zeta = a_{21}\sigma + a_{22}p$$



$$\varepsilon = \frac{1}{K}\sigma + \frac{1}{H^w}p$$

$$\zeta = \frac{1}{H_1^w}\sigma + a_{22}p$$



围压增加，
流体流出

$$\frac{1}{H_1^w} = \frac{1}{H^w}$$

围压作用下，开孔样品
流体含量降低系数

因骨架自身可压缩，
表观体积变化大于挤出的
流体体积

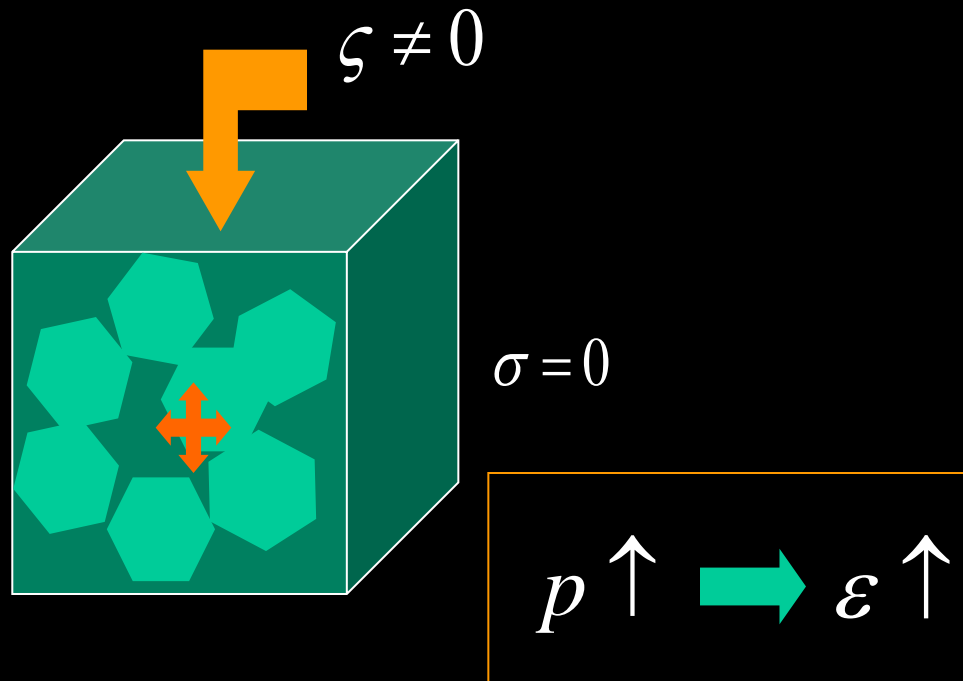
$$\varepsilon = a_{11}\sigma + a_{12}p$$

$$\zeta = a_{21}\sigma + a_{22}p$$



$$\varepsilon = \frac{1}{K}\sigma + \frac{1}{H^w}p$$

$$\zeta = \frac{1}{H_1^w}\sigma + \frac{1}{R}p$$



围压不变，
增加孔压与注入流体同步

$$\frac{1}{R} = \left. \frac{\delta\zeta}{\delta p} \right|_{\sigma} = a_{22}$$

*unconstrained
specific
storage
coefficient*

无约束的比储系数

$$\varepsilon = a_{11}\sigma + a_{12}p$$

$$\zeta = a_{21}\sigma + a_{22}p$$



$$\varepsilon = \frac{1}{K}\sigma + \frac{1}{H^W}p$$

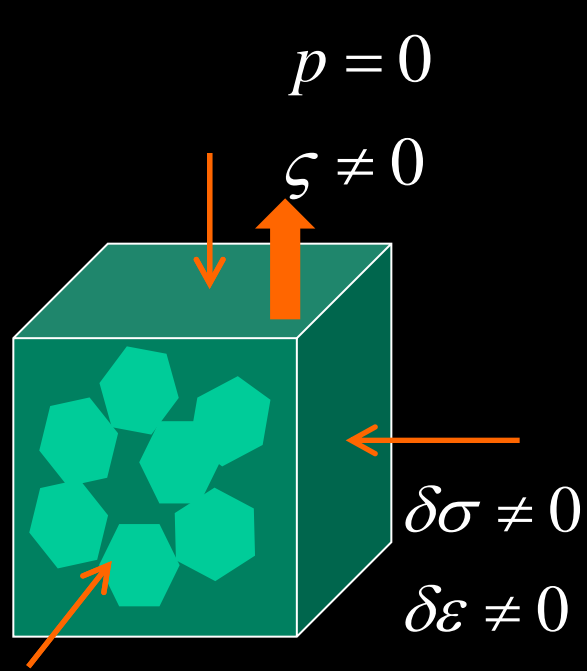
$$\zeta = \frac{1}{H_1^W}\sigma + \frac{1}{R}p$$

$$\varepsilon = \frac{1}{K}\sigma + \frac{\alpha}{K}p$$

$$\zeta = \frac{\alpha}{K}\sigma + \frac{\alpha}{KB}p$$

(2.16-18)

H.F.Wang



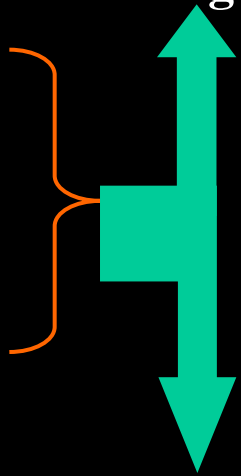
$$\frac{1}{K} = \left. \frac{\delta\varepsilon}{\delta\sigma} \right|_p$$

$$\frac{1}{H^W} = \left. \frac{\delta\zeta}{\delta\sigma} \right|_p$$

$$\frac{K}{H^W} = \left. \frac{\delta\zeta}{\delta\varepsilon} \right|_p$$

$$\left. \frac{\zeta}{\varepsilon} \right|_p = \frac{C}{M} = \alpha$$

$$\frac{K}{H^W} = \frac{C}{M} = \alpha$$



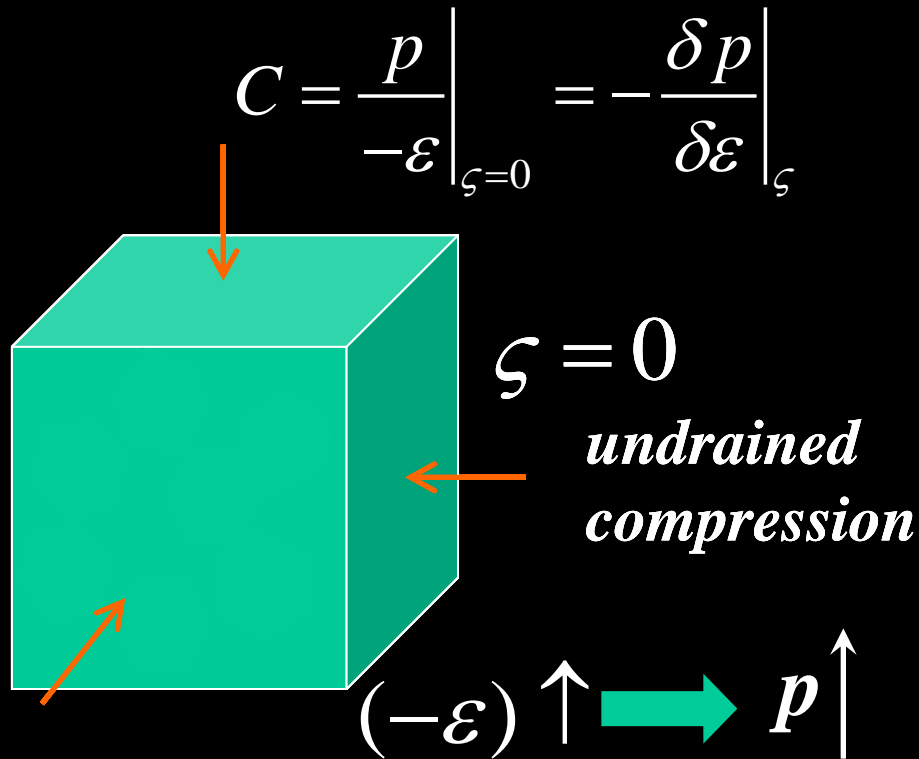
(未封膜, 开孔) 孔压不变。
 增大围压可使体积减小、流体流出。
 α 是流体含量增量与体膨胀量之比

围压作用下, 开孔样品排出的流体体积小于样品减少的表观体积

$$p = -\varepsilon C + M \zeta$$

$$\varepsilon = \nabla \cdot \mathbf{u}, \quad \zeta = -\nabla \cdot \mathbf{w}$$

$$\sigma = K_c \varepsilon - C \zeta$$



C：包膜（不排水）条件下压缩岩样引起的孔隙流体压强增量

$$B = -\frac{\delta p}{\delta \sigma} \Big|_{\zeta}$$

B 不排水条件下, 施加围压而导致的孔隙流体压强增量.

反映孔隙流体与固体骨架在承担围压中的贡献比例

$$B = -\frac{\delta p}{\delta \varepsilon} \Big|_{\zeta} \times \frac{\delta \varepsilon}{\delta \sigma} \Big|_{\zeta} = \frac{C}{K_c} = \frac{\alpha M}{K_c}$$

系数之间的关系

Skempton's [1954] coefficient B and
the Biot and Willis [1957] constant α

$$0 < \alpha < \frac{1}{B}$$

$$\alpha B = \left(1 - \frac{K_d}{K_c} \right) \left\{ \begin{array}{l} B = - \frac{\delta p}{\delta \varepsilon} \Big|_{\zeta} \times \frac{\delta \varepsilon}{\delta \sigma} \Big|_{\zeta} = \frac{C}{K_c} = \frac{\alpha M}{K_c} \\ K_c = \frac{\delta \sigma}{\delta \varepsilon} \Big|_{\zeta} = \alpha C + K_b \end{array} \right.$$

慢纵波对应的渗流位移与固相位移之比

$$\beta = - \frac{1}{B} \left(1 + \frac{4}{3} \frac{G}{K_c} \right)$$

Darcy's Law

$$\dot{\mathbf{w}} = \frac{\kappa}{\eta} \left(\underbrace{(-\nabla p)} + \underbrace{\text{Inertial force on pore fluid}} \right)$$

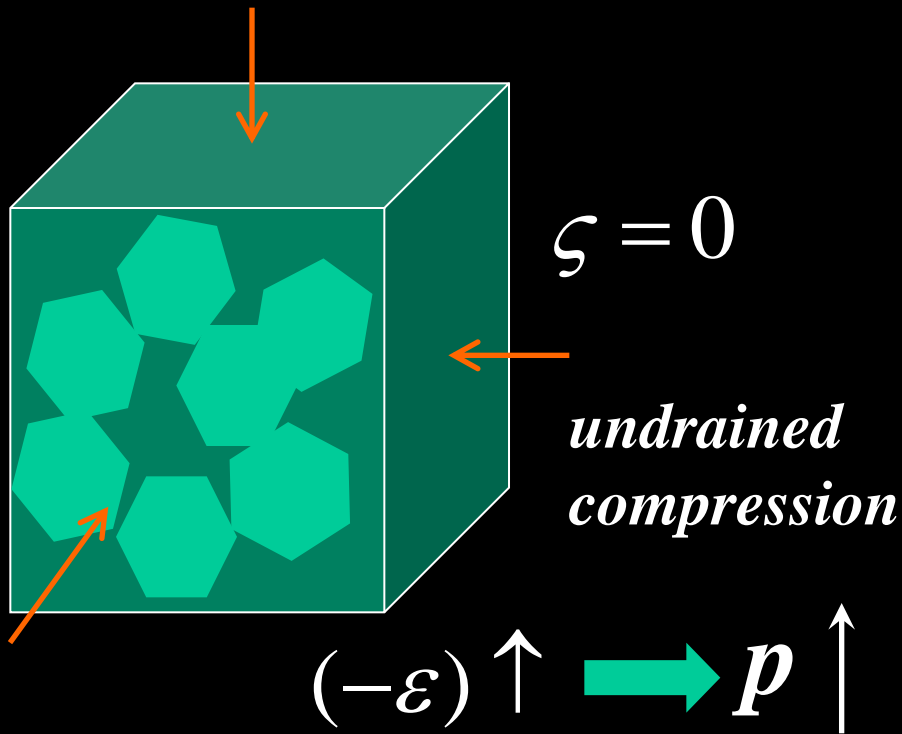
$$-\rho_f \ddot{\mathbf{u}}$$

$$\dot{\mathbf{w}} = \frac{\kappa}{\eta} (-\nabla p + i\omega\rho_f \dot{\mathbf{u}})$$

Biot's equations

$$\varepsilon = \nabla \cdot \mathbf{u}, \quad \zeta = -\nabla \cdot \mathbf{w}$$

$$\sigma = K_c \varepsilon - C \zeta$$



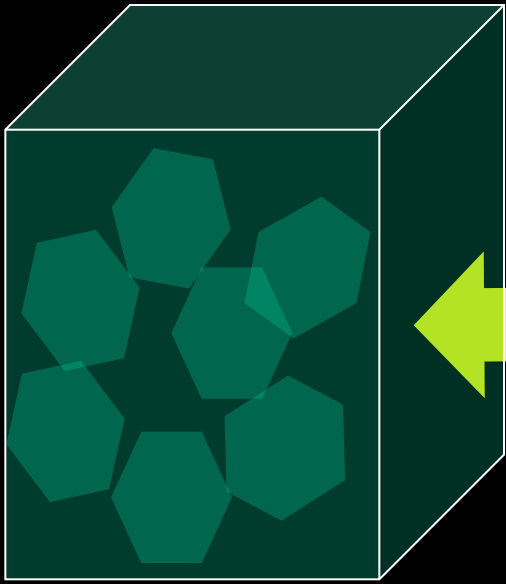
包膜（不排水）条件下，
压缩岩样而增加的孔隙流体压强

$$p = -\varepsilon C + M \zeta$$

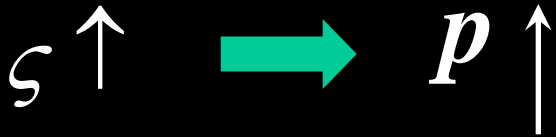
$$C = \left. \frac{p}{-\varepsilon} \right|_{\zeta=0} = - \left. \frac{\delta p}{\delta \varepsilon} \right|_{\zeta}$$

$$K_c = \left. \frac{\delta \sigma}{\delta \varepsilon} \right|_{\zeta} = \alpha C + K_b$$

*input fluid into
a rock sample
within a rigid shell*



$$\varepsilon = 0$$

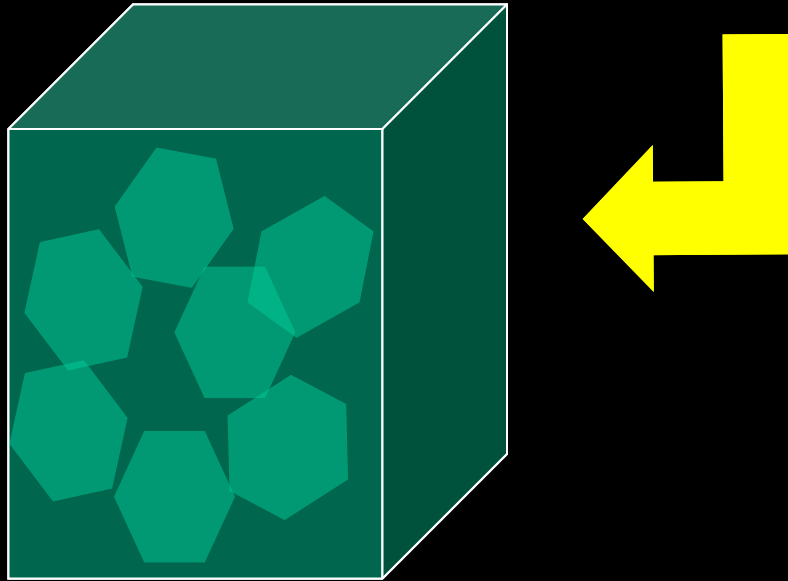


$$p = -\varepsilon C + M \zeta$$

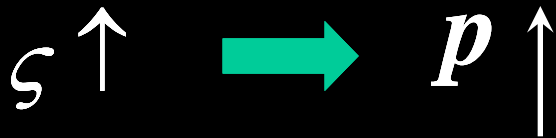
$$M = \left. \frac{p}{\zeta} \right|_{\varepsilon=0} = \left. \frac{\delta p}{\delta \zeta} \right|_{\varepsilon}$$

M : 保持岩样表观体积不变，
因注入流体而增大的流体压强

*input fluid into a
rock sample within a rigid shell*



$$\varepsilon = 0$$



M 保持岩样表观体积不变，
因注入流体而增大的流体压强

$$p = -\varepsilon C + M \zeta$$

$$S_\varepsilon = \frac{1}{M} = \left. \frac{\delta \zeta}{\delta p} \right|_\varepsilon$$

$$S_\sigma = \frac{1}{R} = \left. \frac{\delta \zeta}{\delta p} \right|_\sigma$$

与具有自由边界样品相比，
具有刚性边界的样品（对
注入同样多的流体）呈现
出更大的抵抗力。

$$S_\varepsilon < S_\sigma \longrightarrow M > R$$

$$\frac{1}{M} = \frac{1}{R} - \frac{K}{(H^W)^2}$$

$$S_\varepsilon = S_\sigma - \frac{K}{(H^W)^2}$$

$$p = -\varepsilon C + M\zeta$$

$$S_\varepsilon = \frac{1}{M} = \left. \frac{\delta\zeta}{\delta p} \right|_\varepsilon$$

$$S_\sigma = \frac{1}{R} = \left. \frac{\delta\zeta}{\delta p} \right|_\sigma$$

注入同样多的流体,刚性边界会导致更大的流体压强增量。

$$S_\varepsilon < S_\sigma \quad \longrightarrow \quad M > R$$

Biot's equations

$$\varepsilon = \nabla \cdot \mathbf{u}, \quad \zeta = -\nabla \cdot \mathbf{w}$$

$$\sigma = K_c \varepsilon - C \zeta$$

$$p = -\varepsilon C + M \zeta$$

$$p = 0$$

$$\varepsilon C = M \zeta$$

$$\left. \frac{\zeta}{\varepsilon} \right|_p = \frac{C}{M} = \alpha \leq 1$$

$$C = \left. \frac{p}{-\varepsilon} \right|_{\zeta=0} = - \left. \frac{\delta p}{\delta \varepsilon} \right|_{\zeta}$$


$$C = - \left. \frac{\delta P}{\delta \varepsilon} \right|_{\zeta} \left. \frac{\delta p}{\delta P} \right|_{\zeta} = BK_c$$

$$M = \left. \frac{p}{\zeta} \right|_{\varepsilon=0} = \left. \frac{\delta p}{\delta \zeta} \right|_{\varepsilon}$$

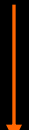
$$M = C / \alpha \geq C$$


$$\frac{1}{K} = \left. \frac{\delta \varepsilon}{\delta \sigma} \right|_p$$

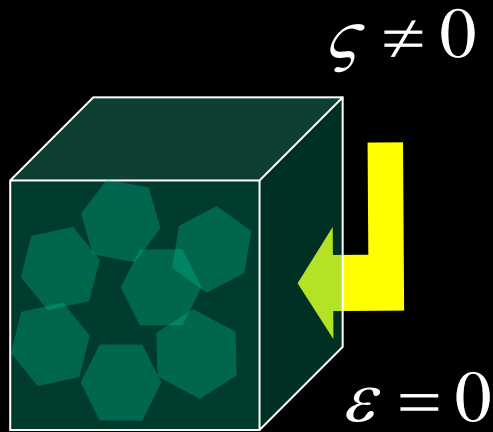
$$\frac{1}{H^W} = \left. \frac{\delta \zeta}{\delta \sigma} \right|_p$$


$$\frac{K}{H^W} = \left. \frac{\delta \zeta}{\delta \varepsilon} \right|_p$$

$$p = -\varepsilon C + M \zeta$$


$$\left. \frac{\zeta}{\varepsilon} \right|_p = \frac{C}{M} = \alpha$$


$$\frac{K}{H^W} = \frac{C}{M} = \alpha$$

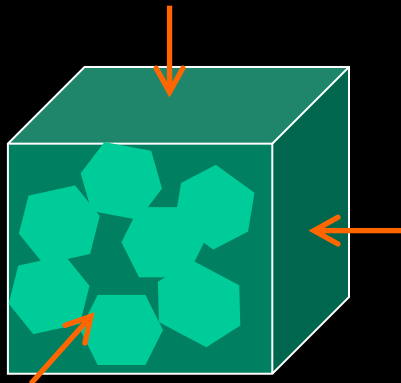


$$M = \frac{p}{\zeta} \Big|_{\varepsilon=0} = \frac{\delta p}{\delta \zeta} \Big|_{\varepsilon}$$

$$M = \frac{1}{\frac{\phi}{K_f} + \frac{\alpha - \phi}{K_s}} = \frac{K_f K_s}{\phi K_s + (\alpha - \phi) K_f}$$

input fluid into a rock sample within a rigid shell

$$K_c = \frac{\delta \sigma}{\delta \varepsilon} \Big|_{\zeta} = \alpha C + K_b$$



$$B = - \frac{\delta p}{\delta \sigma} \Big|_{\zeta} = \frac{\Delta p}{\Delta P} \Big|_{\zeta=0} = \frac{\alpha M}{K_c} = \frac{\alpha M}{\alpha^2 M + K_b}$$

$$B = \frac{\alpha}{\frac{\phi K_c}{K_f} + \frac{(\alpha - \phi) K_c}{K_s}} = \frac{\alpha K_f K_s}{\phi K_s K_c + (\alpha - \phi) K_f K_c}$$

undrained compression

Skempton系数 B

Wang, 2000

$$B = \frac{\Delta p}{\Delta P} \Big|_{\xi=0} = \frac{\alpha M}{K_c} \quad B = \frac{\Delta p}{\Delta P} \Big|_{\xi=0} = \frac{\alpha M}{K_c}, \quad C = K_c B$$



$$\alpha = 1 - \frac{K_b}{K_s}, \quad M = \frac{K_f K_s}{\phi K_s + (\alpha - \phi) K_f},$$

$$K_c = K_b + \alpha^2 M$$

$$K_s, K_f, \phi \implies K_s, K_f, \phi, K_b \implies \alpha, M, K_b \implies K_c$$

$$\mu \implies \rho V_s^2 \quad \frac{4}{3} \mu + K_c = H \implies \rho V_p^2$$

表观体应变的有效应力定理

一块孔隙介质，在两种不同的围压与孔隙组合下分别承受不排水压缩，那么，表观体积改变量只取决于有效应力。

$$P_{ij} = P_{ij} - \alpha p \delta_{ij}$$

特别的，只要有效应力相等，表观体积改变量都等于围压不为零而孔压为零的组合（即排水条件）下的体积改变量。

波动方程向扩散方程的退化

$$\nabla^2 \mathbf{E} = \mu\sigma \frac{\partial \mathbf{E}}{\partial t} + \mu\varepsilon \frac{\partial^2 \mathbf{E}}{\partial t^2}$$

$$k = \omega\sqrt{\mu\varepsilon'}, \quad c = \frac{1}{\sqrt{\mu\varepsilon'}}$$

$\varepsilon\omega \ll \sigma$ ↓ 良导体

$$\nabla^2 \mathbf{E} = \mu\sigma \frac{\partial \mathbf{E}}{\partial t}$$

$$\varepsilon' = \varepsilon + i\frac{\sigma}{\omega}$$

$$\nabla^2 \mathbf{E} + k^2 \mathbf{E} = 0$$

$$k^2 \approx i\omega\mu\sigma = \frac{i\omega}{C_D}, \quad k = \frac{1+i}{\sqrt{2}} \sqrt{\omega\mu\sigma}, \quad Q^{-1} = 2 \frac{\text{Im}(k)}{\text{Re}(k)} = 2$$

$$c^2 \approx \frac{-i\omega}{\mu\sigma} = -i\omega C_D \quad \text{此速度无传播意义}$$

$$\nabla^2 T - \frac{1}{C_D} \frac{\partial T}{\partial t} = 0,$$

热流密度与温度梯度呈正比

$$C_D = \frac{k}{\rho c}, \quad q = -k\nabla T$$

扩散率=热传导系数与（密度与比热的积）之比

孔隙介质慢纵波

$$V_{P2}^2 \approx -i\omega C_D$$

$$C_D = \frac{\kappa K_f}{\eta\phi} \left(1 + \frac{K_f}{\phi \left(K_b + \frac{4}{3}N \right)} \left\{ 1 + \frac{1}{K_s} \left[\frac{4}{3}N \left(1 - \frac{K_b}{K_s} \right) - K_b - \phi \left(K_b + \frac{4}{3}N \right) \right] \right\} \right)^{-1}$$

Chandler and Johnson, 1981

谢谢！

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